



## Optyka nanostruktur

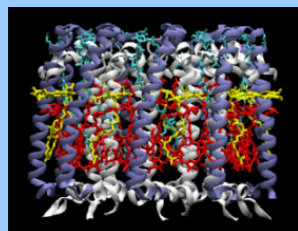
Sebastian Maćkowski

Instytut Fizyki

Uniwersytet Mikołaja Kopernika

Adres poczty elektronicznej: mackowski@fizyka.umk.pl

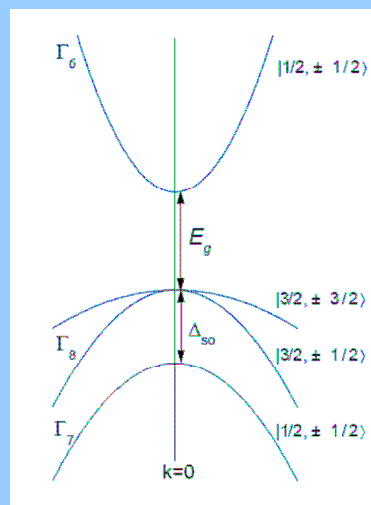
Biuro: 365, telefon: 611-3250



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## Struktura pasmowa



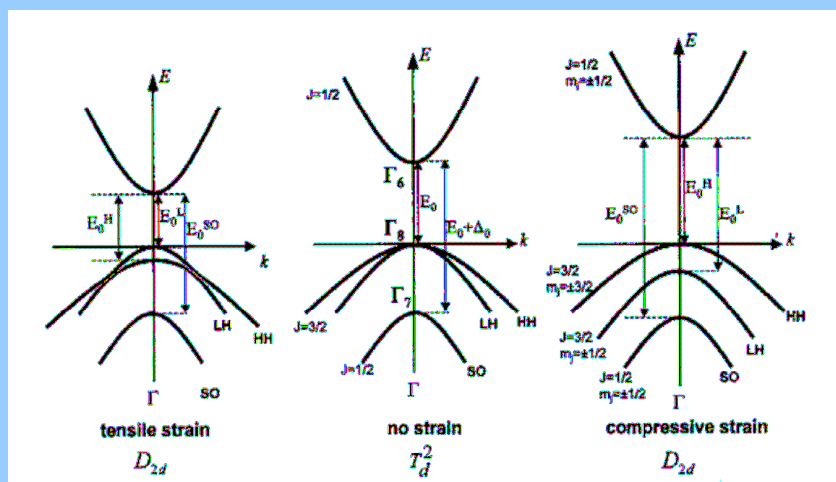
przejścia optycznie dozwolone:  
 $\Delta s = \pm 1$  – bright exciton

przejścia wzbronione  
 $\Delta s = \pm 2$  – dark exciton

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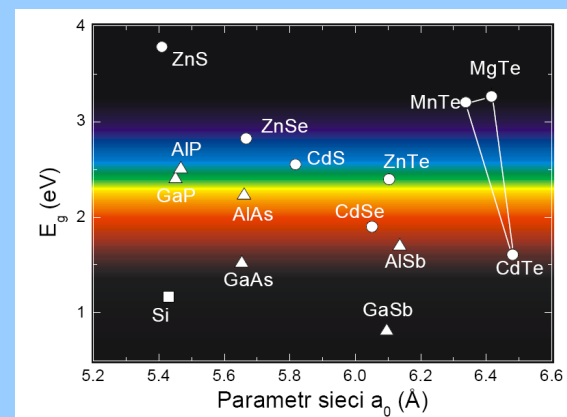
## Wpływ naprężeń



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## Stała sieci vs. energia przerwy



dla większości układów występuje  
niedopasowanie sieciowe!

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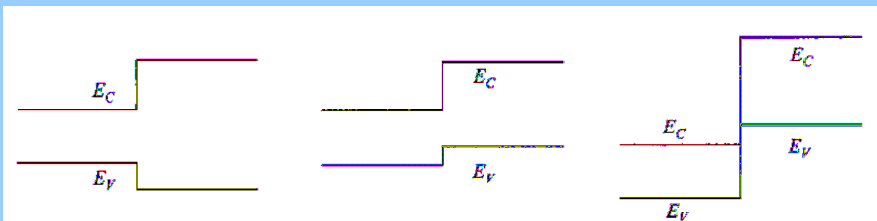


## Heterozłącze

typ I

typ II

typ III



## Studnia kwantowa

AlGaAs GaAs AlGaAs



dla nieskończonych barier

$$E_n = \frac{\hbar^2 \pi^2 n^2}{mL^2}$$

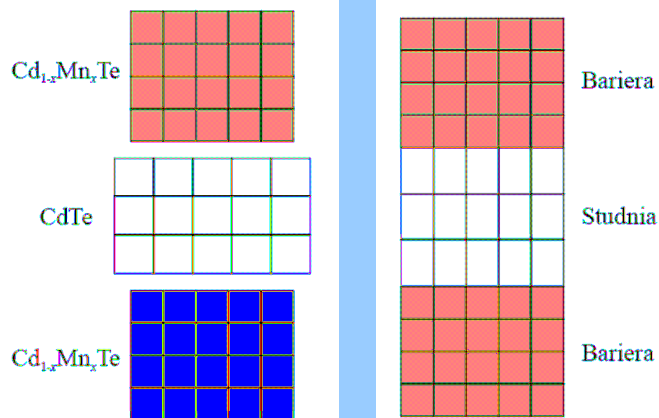
dla skończonych barier

$$V_h(z_h) = \begin{cases} 0, & |z_h| < \frac{1}{2}L \\ V_h, & |z_h| > \frac{1}{2}L \end{cases}$$

$$V_e(z_e) = \begin{cases} 0, & |z_e| < \frac{1}{2}L \\ V_e, & |z_e| > \frac{1}{2}L \end{cases}$$



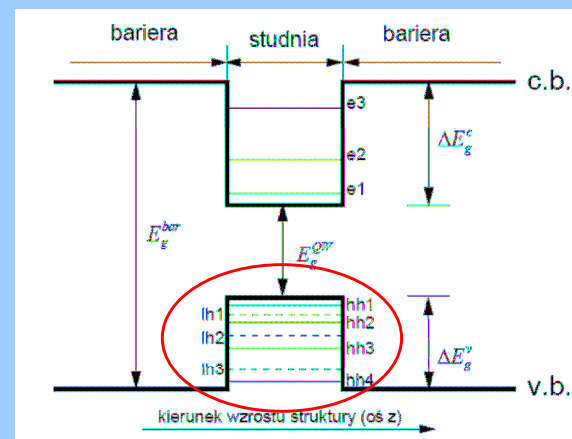
## Wzrost pseudomorficzny



napężenie ściskające



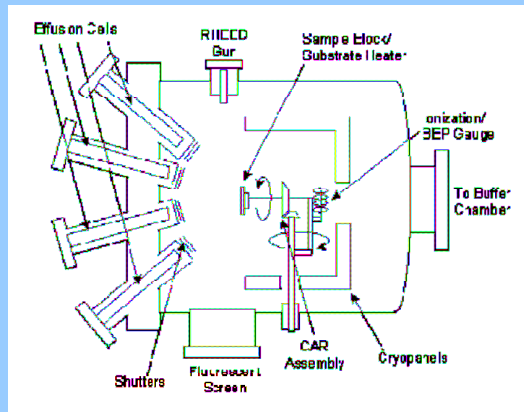
## Zniesienie degeneracji pasm





## Epitaksja

epitaksja z wiązek molekularnych (MBE)



$p \sim 10^{-9}$ - $10^{-10}$  Torr  
1 Torr – 133 Pa

charakteryzacja  
*in-situ*

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## Urządzenie MBE



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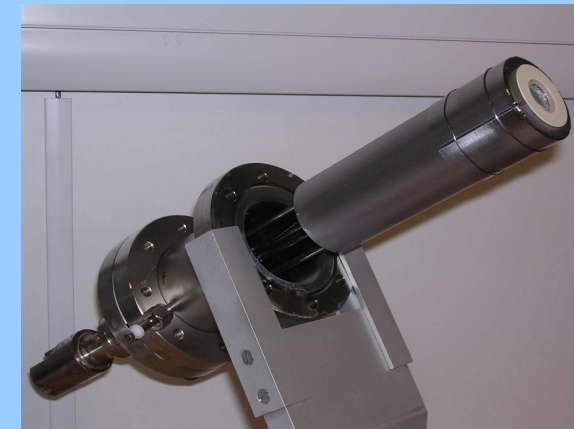
## Możliwości MBE

- wytwarzanie struktur wielowarstwowych złożonych ze związków o różnych strukturach elektronowych, np. GaAs – AlGaAs (band-gap engineering = inżynieria przerwy energetycznej)
- wysoki stopień czystości chemicznej osadzanych warstw
- dobra kontrola ostrości interfejsu (zmiana składu możliwa w jednej warstwie atomowej), możliwość otrzymywania pojedynczej warstwy atomowej (1 ML = 1 monolayer)
- otrzymywanie supersieci (np. GaAs-AlGaAs-GaAs- itd.) o okresie od kilku Å.
- szerokie zakresy domieszkania, np. dla (Al)GaAs, często w zakresie niedostępnym przy wykorzystaniu klasycznych metod wzrostu

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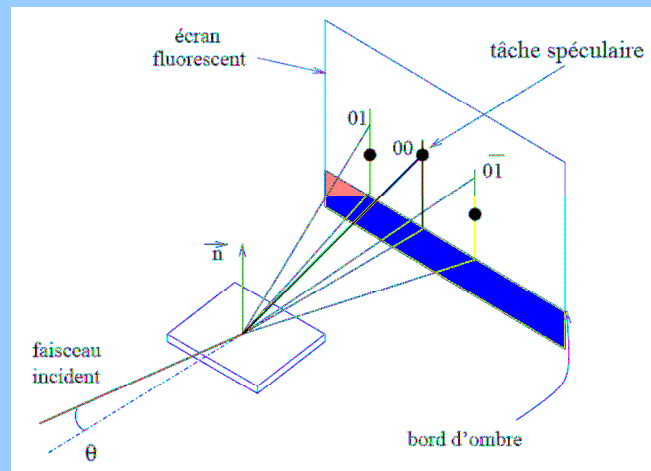
## Komórka efuzyjna



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## RHEED

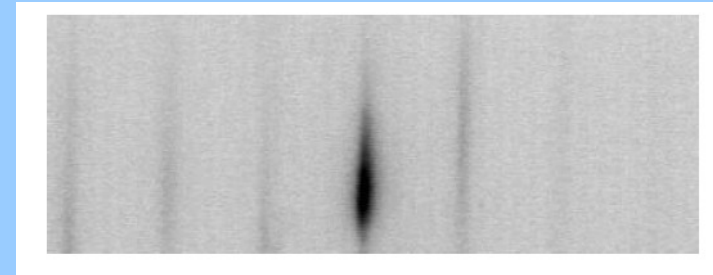


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## RHEED

warstwa CdTe na podłożu ZnTe

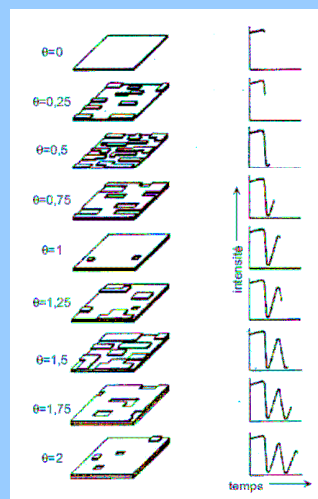
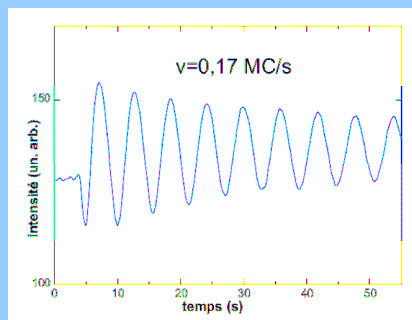


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## Oscylacje RHEED

oscylacje natężenia pozwalają określić szybkość wzrostu



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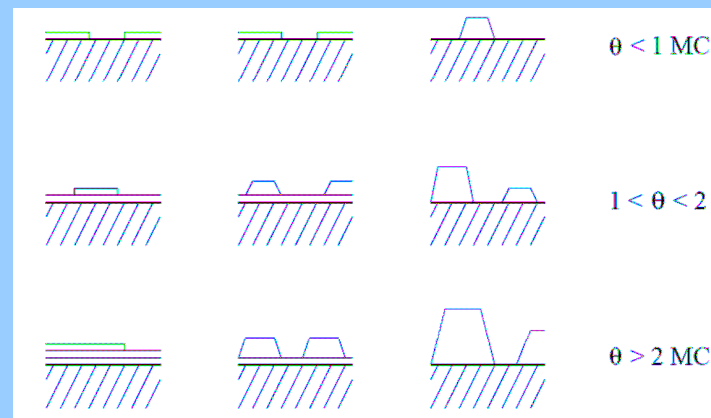


## Mody wzrostu

Frank van der Merwe

Stranski-Krastanov

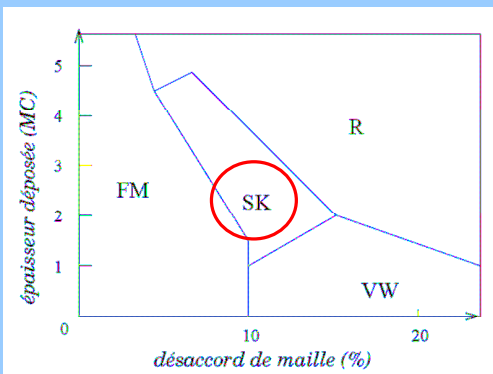
Volmer-Weber



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## Diagram fazowy



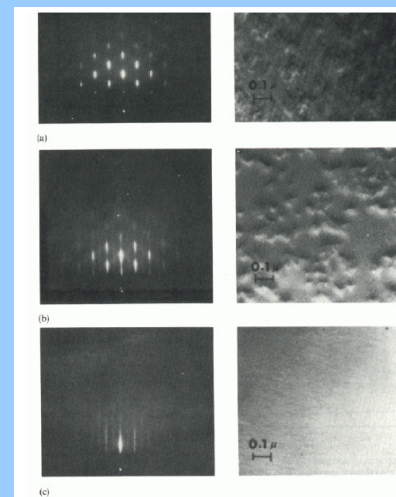
Composés	Désaccord de maille (%)
<i>InAs/GaAs</i>	7,2
<i>CdSe/ZnSe</i>	6,8
<i>GaN/AlN</i>	2,4(c); 2,7(h)
<i>Si<sub>0,7</sub>Ge<sub>0,3</sub>/Si</i>	1,2
<i>CdTe/ZnTe</i>	6,2

tak powstają kropki kwantowe

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## Obraz RHEED



wzrost  
wyspowy

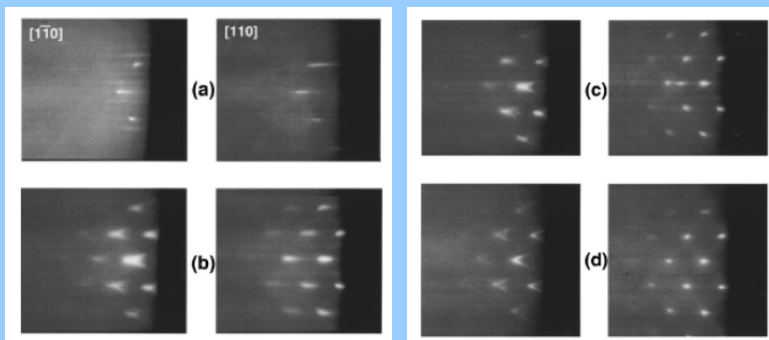
Frank van  
der Merwe

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## Obraz RHEED

RHEED w przypadku wzrostu warstw niedopasowanych sieciowo



wzrost InAs na GaAs(001)  
przy ok. 2 ML wzrost przechodzi w mod Stransky-Krastanov'a

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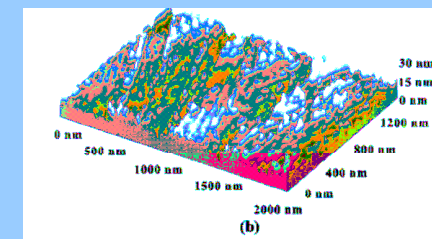
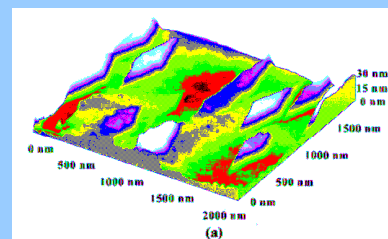


## AFM

AFM – mikroskopia sił atomowych jako metoda charakteryzacji  
powierzchni w ultra wysokiej próżni

3 ML CdTe na ZnTe

6 ML CdTe na ZnTe



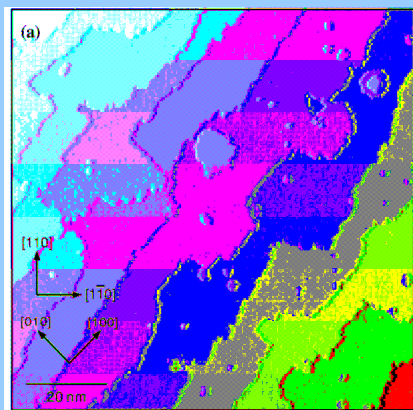
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## STM

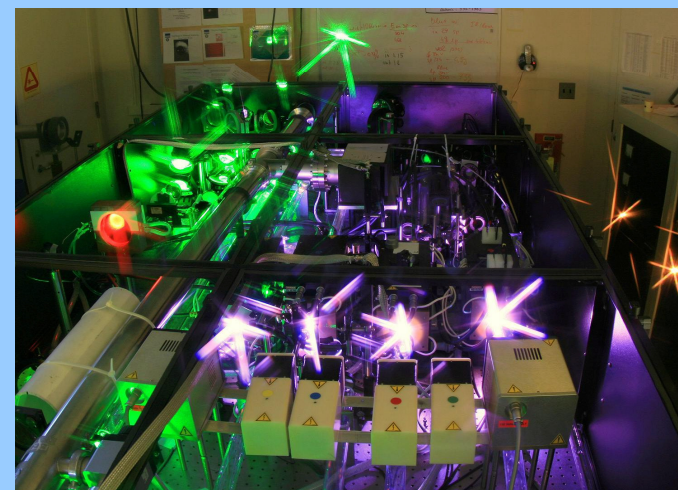
STM – skaningowa mikroskopia tunelowa jako metoda charakteryzacji powierzchni w ultra wysokiej próżni



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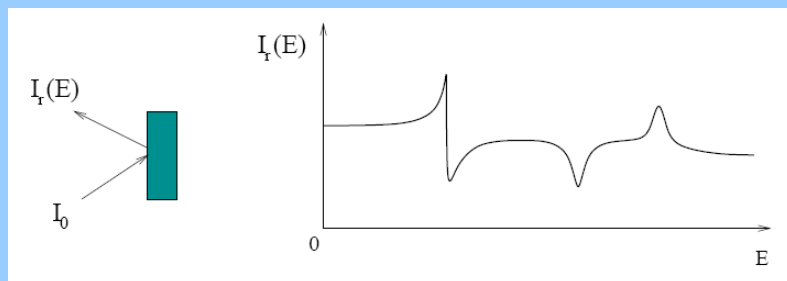
## Spektroskopia optyczna



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## Odbicie

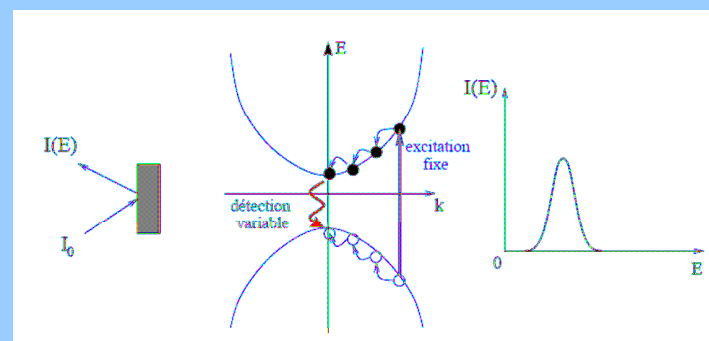


energie ekscytonów – pomiar prostszy niż absorpcja

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## Luminescencja

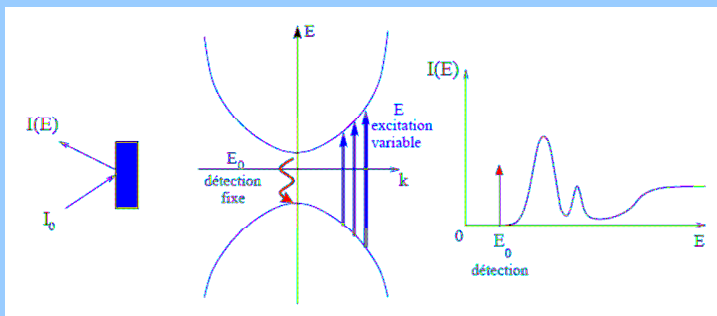


dynamika ekscytonów, mechanizmy poszerzenia

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## Wzbudzenie luminescencji

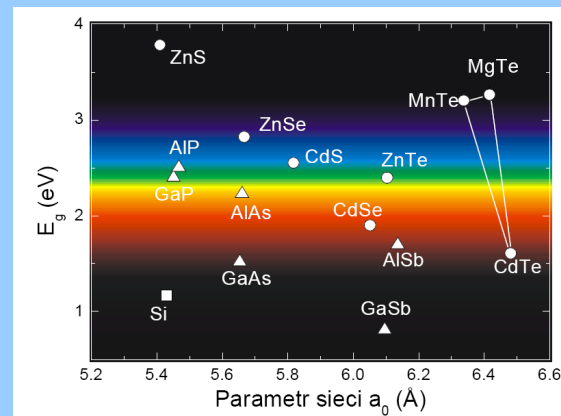


wzbudzenie ekscytonów,  
mechanizmy relaksacji

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## Pierwsze studnie



GaAs-AlGaAs – dla 30% Al brak niedopasowania sieciowego

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## Pierwsze studnie

### Quantum States of Confined Carriers in Very Thin $\text{Al}_x\text{Ga}_{1-x}\text{As-GaAs-Al}_x\text{Ga}_{1-x}\text{As}$ Heterostructures

R. Dingle, W. Wiegmann, and C. H. Henry  
*Bell Laboratories, Murray Hill, New Jersey 07974*  
(Received 24 June 1974)

Quantum levels associated with the confinement of carriers in very thin, molecular-beam-grown  $\text{Al}_x\text{Ga}_{1-x}\text{As-GaAs-Al}_x\text{Ga}_{1-x}\text{As}$  heterostructures result in pronounced structure in the GaAs optical absorption spectrum. Up to eight resolved exciton transitions, associated with different bound-electron and bound-hole states, have been observed. The heterostructure behaves as a simple rectangular potential well with a depth of  $\approx 0.88\Delta E_g$  for confining electrons and  $\approx 0.12\Delta E_g$  for confining holes, where  $\Delta E_g$  is the difference in the semiconductor energy gaps.

testowanie mechaniki kwantowej!

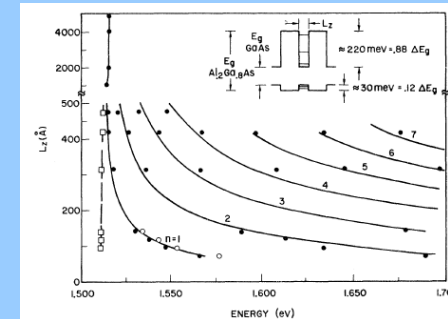
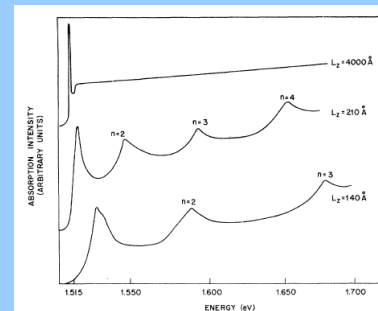
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## Pierwsze studnie

### Quantum States of Confined Carriers in Very Thin $\text{Al}_x\text{Ga}_{1-x}\text{As-GaAs-Al}_x\text{Ga}_{1-x}\text{As}$ Heterostructures

R. Dingle, W. Wiegmann, and C. H. Henry  
*Bell Laboratories, Murray Hill, New Jersey 07974*  
(Received 24 June 1974)



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## Luminescencja

PHYSICAL REVIEW B VOLUME 22, NUMBER 2 15 JULY 1980

### Luminescence studies of optically pumped quantum wells in GaAs-Al<sub>x</sub>Ga<sub>1-x</sub>As multilayer structures

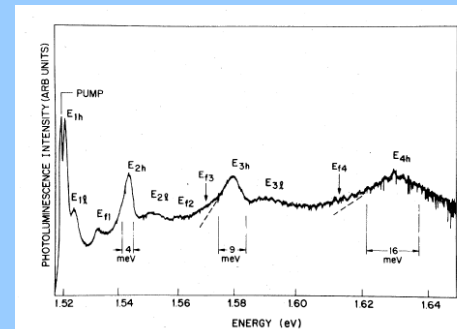
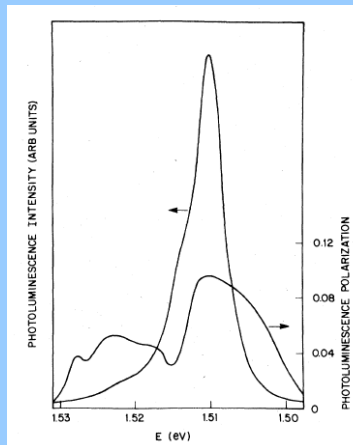
R. C. Miller, D. A. Kleinman, W. A. Nordland, Jr., and A. C. Gossard  
Bell Laboratories, Murray Hill, New Jersey 07974  
(Received 9 January 1980)

Various properties derivable from the photoluminescence of a multilayer structure consisting of many thin GaAs layers (wells) separated by thin Al<sub>x</sub>Ga<sub>1-x</sub>As layers (barriers) have been determined as a function of temperature from 7 K to greater than 150 K. These measurements, which were made utilizing both fixed- and tunable-frequency laser sources, include the following: the circular photoluminescence polarization  $\rho$  generated by circularly polarized optical excitation, the magnetic field dependence of  $\rho$  in the Voigt geometry (Hanle effect) to determine the electron lifetime  $\tau$  and spin-relaxation time  $\tau_s$ , the excitation spectrum for the peak of the photoluminescence which shows thirteen exciton quantum transitions characteristic of multilayer samples, the photoluminescence spectra which exhibit five transitions in emission, the excitation intensity  $I_p$  dependence of the integrated photoluminescence intensity, and  $\tau$  and  $\tau_s$  as a function of  $I_p$ . The temperature dependence of  $\tau_s$  suggests that there is a center with a binding energy of about 5 meV that relaxes the electron spins at low temperatures. In the high-temperature region,  $T \geq 100$  K,  $\tau_s$  is close to that observed by others for lightly doped p-type bulk GaAs and is close to estimates of  $\tau_s$  for the D'yakonov-Perel' spin-relaxation mechanism. The absence of optical alignment with linearly polarized pumping and the characteristics of the magnetic field dependence of  $\rho$  both argue against significant hole polarization. The linear dependence of the integrated photoluminescence intensity on  $I_p$  leads to a simple model for the radiative lifetime whose temperature dependence is found to be consistent with expectations. Also, these data are used to estimate the nonradiative lifetime  $\tau_{nr} = 5 \times 10^{-10}$  sec which varies little with  $T$  and the radiative lifetime  $\tau_r(T) = 3 \times 10^{-10}$  sec at 7 K, which varies  $\sim T^{3/2}$ .

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## Luminescencja

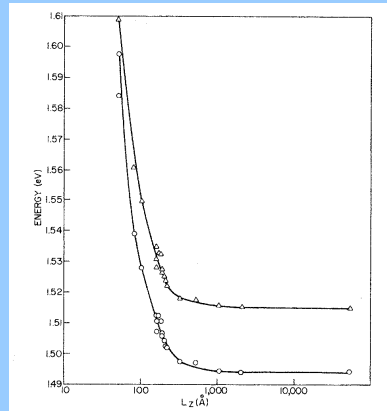
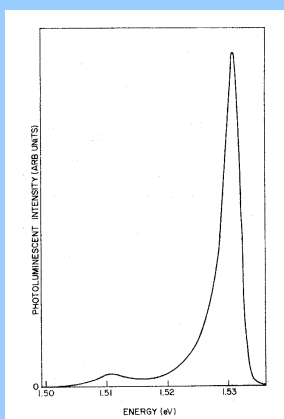


szerokość linii rezonansowych

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## Luminescencja

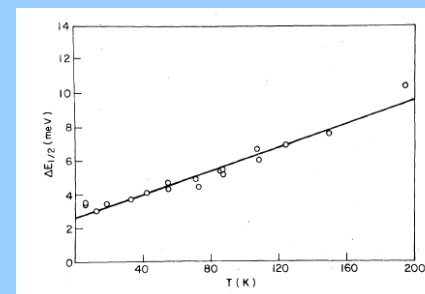


zależność energii emisji ekscytonu od szerokości studni kwantowej

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## Luminescencja



poszerzenie linii luminescencyjnej związane z rozpraszaniem na fononach akustycznych i optycznych

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## Biekscyton

PHYSICAL REVIEW B

VOLUME 25, NUMBER 10

15 MAY 1982

### Biexcitons in GaAs quantum wells

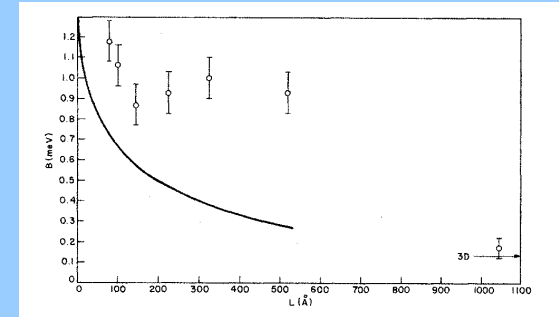
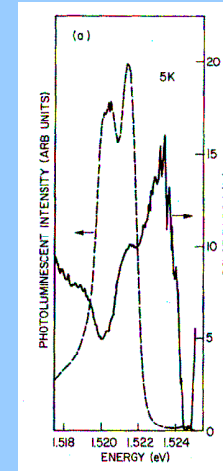
R. C. Miller, D. A. Kleinman, A. C. Gossard, and O. Munteanu  
*Bell Laboratories, Murray Hill, New Jersey 07974*  
(Received 24 March 1982)

Careful examination of the exciton edge of the photoluminescence from a number of high-quality multiple-GaAs-quantum-well samples grown by molecular-beam epitaxy reveals at low temperatures a double peak whose splitting of approximately 1 meV decreases somewhat with increasing GaAs well width  $L$ . The higher-energy peak is due to the  $n=1$  free-heavy-hole-exciton transition while the excitation intensity, temperature, and polarization dependences of the lower-energy peak suggest that it is due to biexcitons with a binding energy  $B$  of about 1 meV. In support of the biexciton interpretation a theoretical calculation of  $B(L)$  is presented. This calculation gives two-dimensional biexciton binding energies more than an order of magnitude larger than the three-dimensional calculated values.

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## Biekscyton



zależność energii wiązania biekscytonu od szerokości studni kwantowej

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## Biekscyton

PHYSICAL REVIEW B

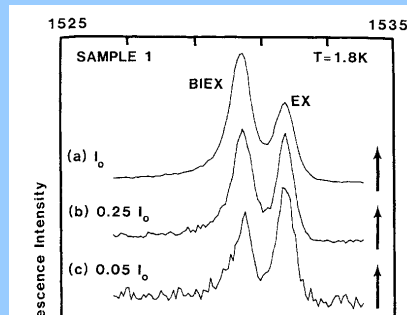
VOLUME 38, NUMBER 5

15 AUGUST 1988-1

### Optical investigation of biexcitons and bound excitons in GaAs quantum wells

S. Charbonneau, T. Steiner, and M. L. W. Thewalt  
*Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada V5A 1S6*

Emil S. Koteles, J. Y. Chi, and B. Elman  
*General Telephone and Electronics Laboratories Incorporated, 40 Sylvan Road, Waltham, Massachusetts 02254*  
(Received 26 February 1988)



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## Tunelowanie

PHYSICAL REVIEW B

VOLUME 41, NUMBER 17

15 JUNE 1990-1

### Resonant-tunneling transfer times between asymmetric GaAs/Al<sub>0.35</sub>Ga<sub>0.65</sub>As double quantum wells

M. G. W. Alexander, M. Nido,\* and W. W. Rühle  
*Max-Planck-Institut für Festkörperforschung, D7000 Stuttgart 80, Federal Republic of Germany*

K. Köhler  
*Fraunhofer-Institut für Angewandte Festkörperphysik, D7800 Freiburg, Federal Republic of Germany*  
(Received 26 January 1990; revised manuscript received 23 March 1990)

Electron tunneling through the barrier in asymmetric double-quantum-well structures is investigated by time-resolved picosecond luminescence spectroscopy. Change from nonresonant to resonant tunneling is achieved with a perpendicular electric field. Energetic alignment of electron subbands in the two wells strongly enhances tunneling transfer rates. The resonant transfer times decrease strongly with barrier thickness. The wells are coupled at resonance by energy-conserving scattering processes between states localized in a single well. The buildup of delocalized coherent states at resonance would lead to much shorter transfer times.

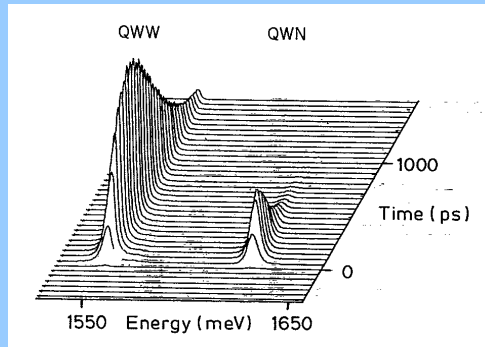
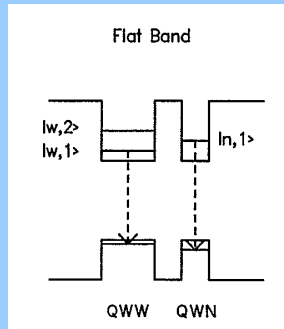
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## Podwójna studnia kwantowa

struktura pasmowa

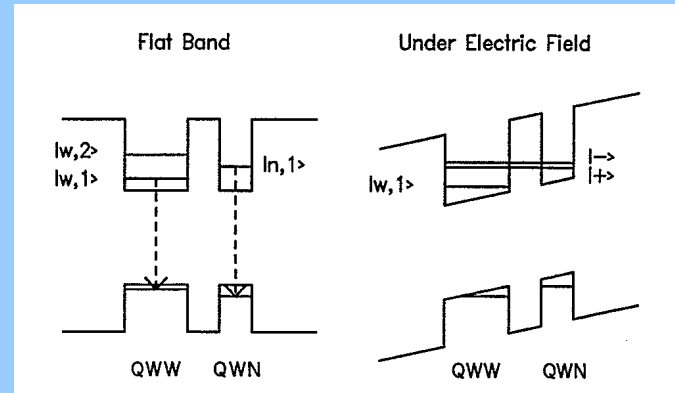
dynamika luminescencji



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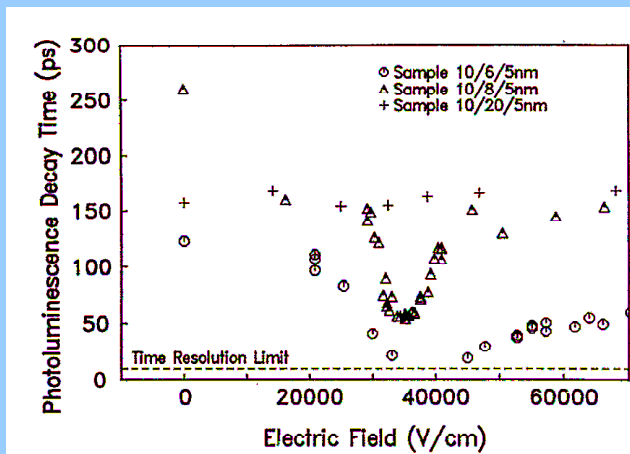
## Podwójna studnia kwantowa



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## Tunelowanie rezonansowe



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## Ekscyton skośny

PHYSICAL REVIEW B VOLUME 38, NUMBER 9 15 SEPTEMBER 1988-II

### Brief Reports

*Brief Reports are short papers which report on completed research which, while meeting the usual Physical Review standards of scientific quality, does not warrant a regular article. (Addenda to papers previously published in the Physical Review by the same authors are included in Brief Reports.) A Brief Report may be no longer than 3½ printed pages and must be accompanied by an abstract. The same publication schedule as for regular articles is followed, and page proofs are sent to authors.*

#### Transformation of spatially direct to spatially indirect excitons in coupled double quantum wells

S. Charbonneau and M. L. W. Thewalt  
*Department of Physics, Simon Fraser University, Burnaby, British Columbia, Canada V5A 1S6*

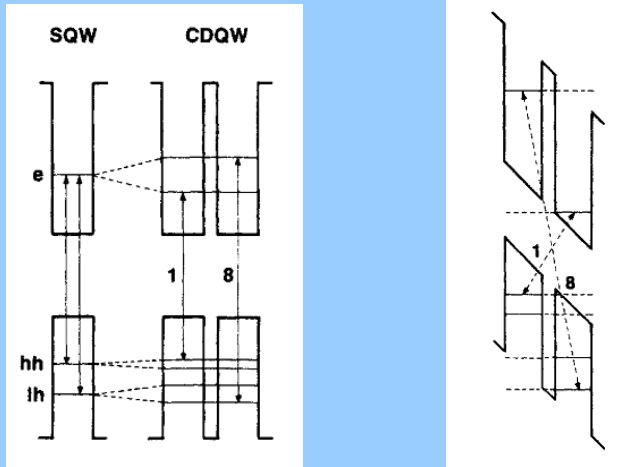
Emil S. Koteles and B. Elman  
*GTE Laboratories Incorporated, 40 Sylvan Road, Waltham, Massachusetts 02254*  
(Received 12 April 1988)

We report a detailed experimental study of the influence of electric fields on exciton lifetimes in a GaAs/Al<sub>x</sub>Ga<sub>1-x</sub>As coupled-double-quantum-well structure. The energy of the lowest-lying photoluminescent exciton peak was observed to decrease and its decay time was found to increase by over an order of magnitude as the internal electric field was increased. The increase in the decay time is attributed to the change of the exciton transition from spatially direct (i.e., taking place within a single well) to spatially indirect (connecting an electron level in one well with a hole level in the adjacent well).

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# Ekscyton skośny



# Ekscyton skośny

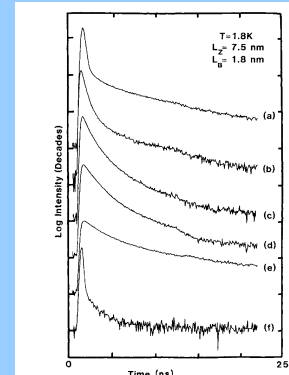
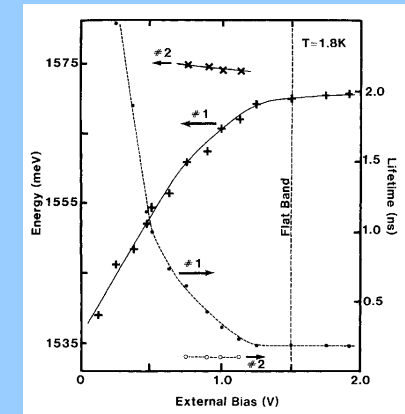


FIG. 3. The luminescence decay of exciton peaks PL [curves (a)-(e) correspond to transition 1 and curve (f) corresponds to transition 2] under various bias voltages. Each of the tick marks on the vertical scale represents one decade of intensity. The transient curves represent decays having time constants of (a)  $V = 1.5$  V,  $\tau = 0.19$  ns; (b)  $V = 1.0$  V,  $\tau = 0.31$  ns; (c)  $V = 0.75$  V,  $\tau = 0.61$  ns; (d)  $V = 0.50$  V,  $\tau = 1.0$  ns; (e)  $V = 0.375$  V,  $\tau = 1.91$  ns; (f)  $V = 1.0$  V,  $\tau = 0.10$  ns (transition no. 2).



# Efekt Starka

VOLUME 55, NUMBER 23      PHYSICAL REVIEW LETTERS      2 DECEMBER 1985

## Lifetime Enhancement of Two-Dimensional Excitons by the Quantum-Confined Stark Effect

H.-J. Pollard, L. Schultheis, and J. Kuhl  
Max-Planck-Institut für Festkörperforschung, D-7000 Stuttgart 80, Federal Republic of Germany

E. O. Göbel  
Fachbereich Physik, Philipps-Universität, D-3550 Marburg, Federal Republic of Germany

and

C. W. Tu  
AT&T Bell Laboratories, Murray Hill, New Jersey 07974  
(Received 9 September 1985)

We report on picosecond luminescence studies of GaAs/AlGaAs quantum wells in the regime of the quantum-confined Stark effect. A drastic increase of the recombination lifetime is accompanied by a Stark shift of the photoluminescence of the lowest free exciton for electric fields perpendicular to the quantum-well layers. A consistent picture of the quantum-confined Stark effect is presented.



# Efekt Starka

